

# Universal scaling and the essential singularity at the abrupt Ising transition

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# Outline

- ▶ Renormalization and the Ising model

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- ▶ Metastability and complex free energy

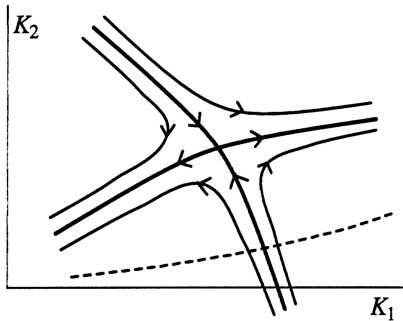
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- ▶ Closed-form results for the 2D Ising susceptibility

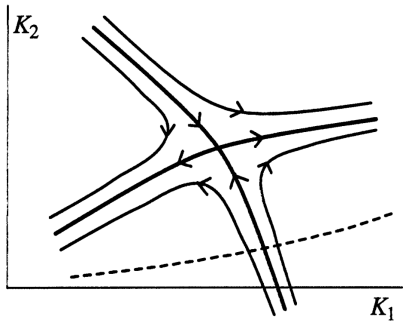
# Renormalization and the Ising Model



From *Scaling and Renormalization in Statistical Physics* by John Cardy

RG analytically maps system space onto itself.

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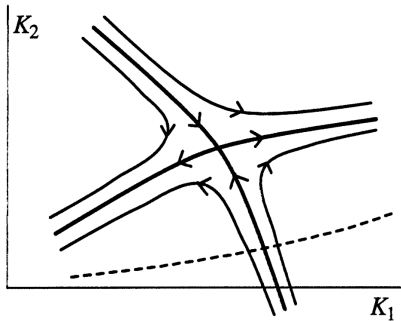


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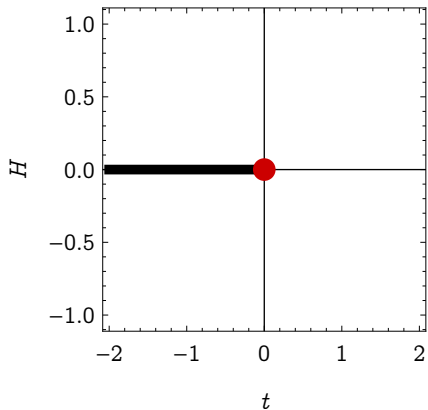
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Fixed points correspond to phases, criticality.

Nonanalytic behavior is preserved by RG.

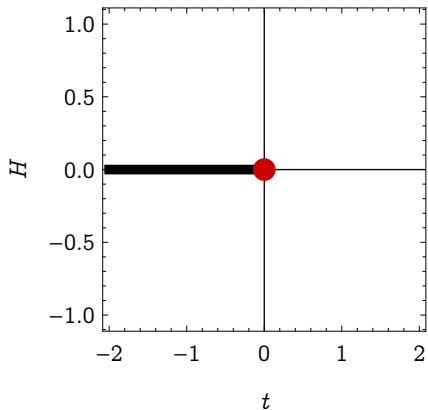


# Renormalization and the Ising Model



Ising critical point has power laws, logarithms in thermodynamic variables.

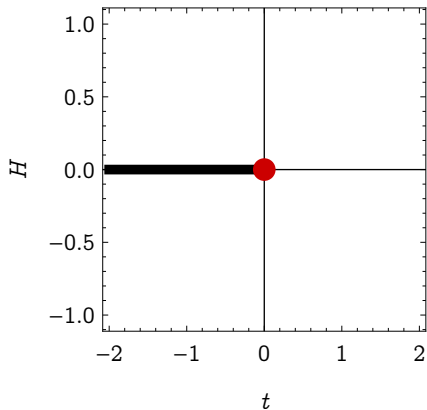
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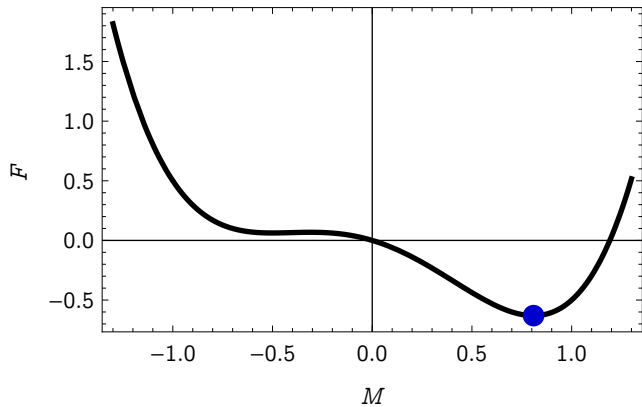
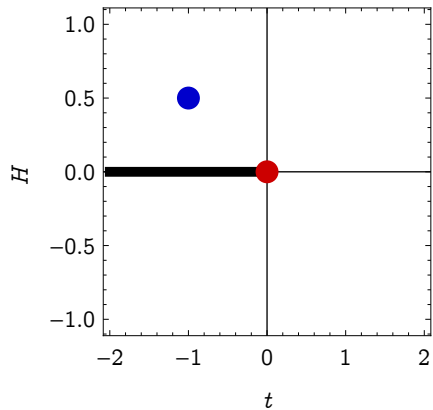


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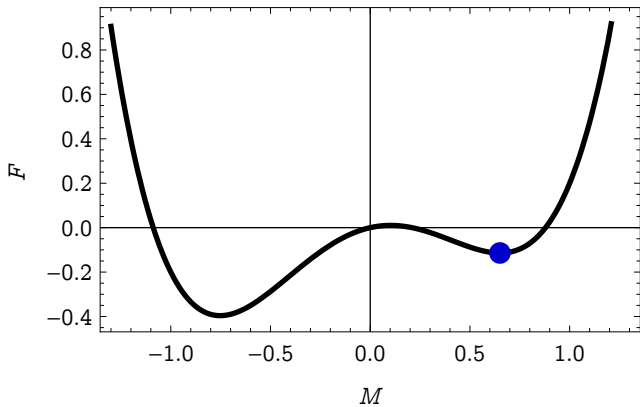
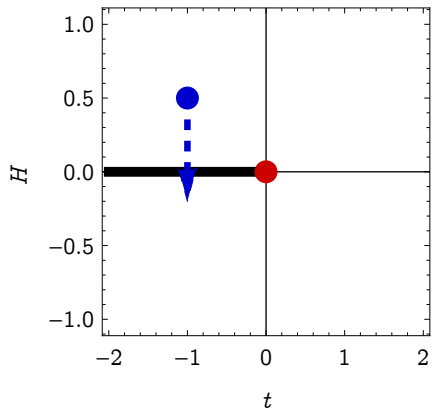
Connected to line of abrupt transitions.

We've identified nonanalytic behavior along the abrupt transition line.

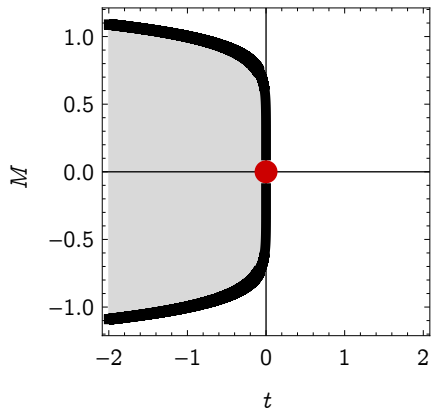
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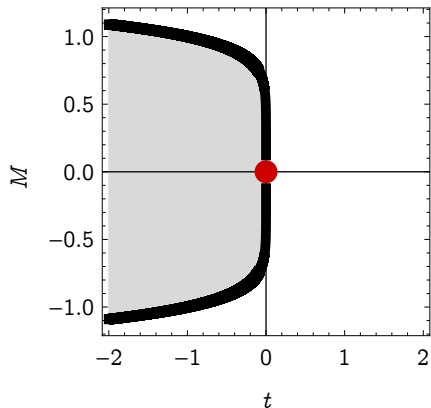


## Metastability & Complex Free Energy



Thermodynamics can be continued into metastable phase.

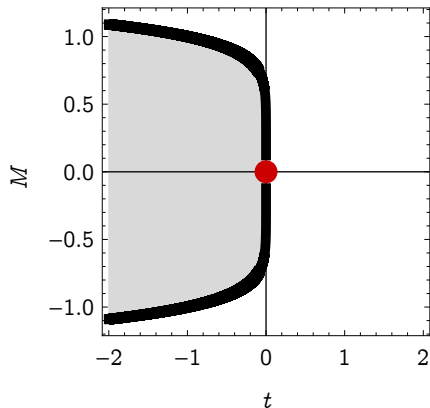
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Ising metastable decay somewhat well studied (Günther 1980, Houghton 1980)



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Metastable phase is stable to domains smaller than

$$N_{\text{crit}} = \left( \frac{MH}{\sigma \Sigma} \right)^{-1/(\sigma-1)}$$

but larger will grow to occupy the entire system, decay to stable phase.

# Metastability & Complex Free Energy

Formation of critical domain has energy cost

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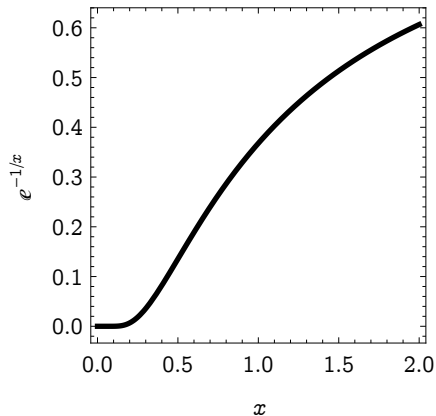
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Imaginary free energy is therefore

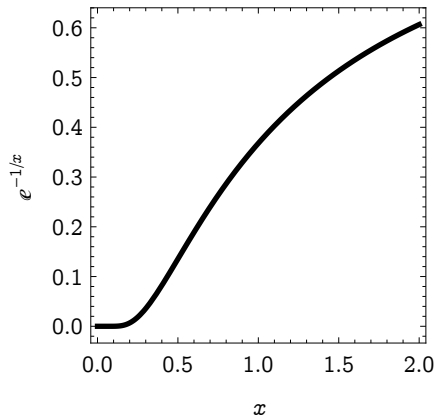
$$\text{Im } F \sim \Gamma \sim P_{\text{crit}} \sim e^{-\beta \Delta F_{\text{crit}}} = e^{-\beta (\Sigma / (MH)^\sigma)^{1/(1-\sigma)}}$$

## Metastability & Complex Free Energy



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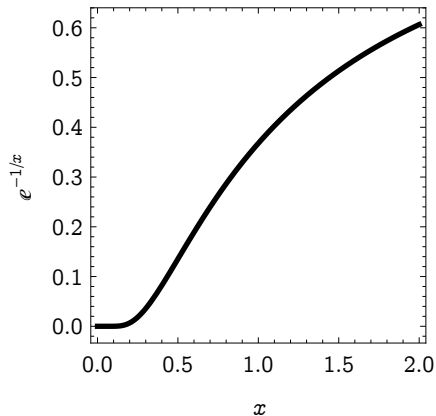
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Near critical point,  $\sigma = 1 - \frac{1}{d}$ , and

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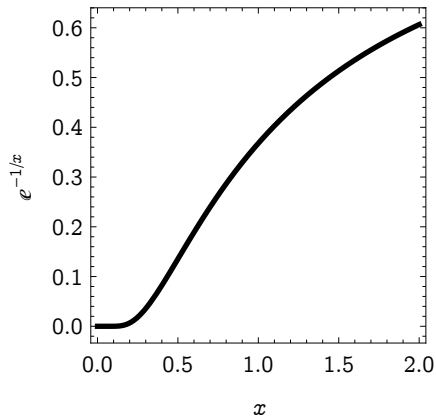
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Nonanalytic behavior is universal!

Can directly observe by measuring metastable decay rate, but what else?

# Analytic Constrains on the Stable Free Energy

# The Metastable Ising Model

Near the Ising critical point,  $\sigma = 1 - \frac{1}{d}$  and

$$M = t^\beta \mathcal{M}(h/t^{\beta\delta})$$

$$\Sigma = t^\mu \mathcal{S}(h/t^{\beta\delta})$$

with  $\mathcal{M}(0)$  and  $\mathcal{S}(0)$  nonzero and finite.

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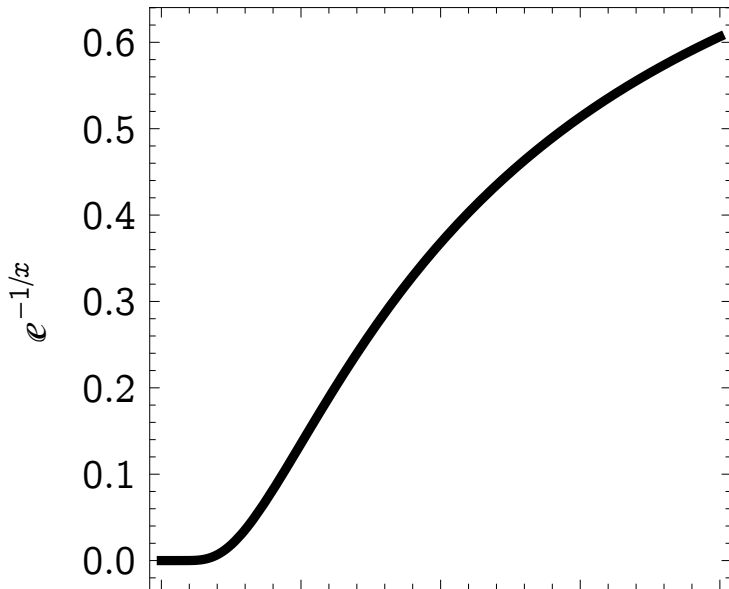
Therefore,

$$\Delta F_{\text{crit}} \sim \Sigma \left( \frac{MH}{\Sigma} \right)^{-(d-1)} = X^{-(d-1)} \mathcal{F}(X)$$

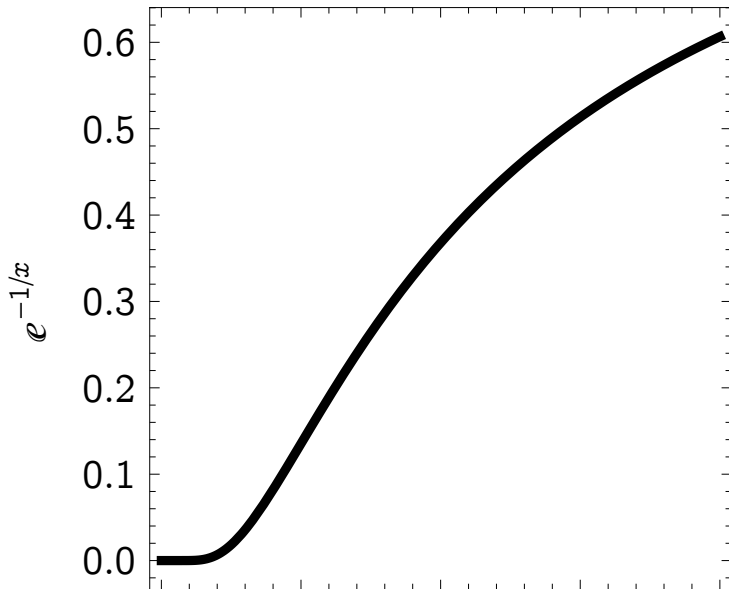
for  $X = h/t^{\beta\delta}$ , and

$$\text{Im } F = t^{2-\alpha} \mathcal{I}(X) e^{-\beta/X^{(d-1)}}$$

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Only predictive for high moments of  $F$ , or

$$f_n = \frac{1}{\pi} \int_{-\infty}^0 \frac{\operatorname{Im} F(X')}{X'^{n+1}} dX'$$

for  $F = \sum f_n X^n$ .

# The Essential Singularity

Results from field theory indicate that  $\mathcal{I}(X) \propto X + \mathcal{O}(X^2)$  for  $d = 2$ , so that

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Not a convergent series—the real part of  $F$  for  $H > 0$  is also nonanalytic!

## The Essential Singularity

In two dimensions, the Cauchy integral does not converge, normalize with  $\lambda$ ,

$$F(X | \lambda) = \frac{1}{\pi} \int_{-\infty}^0 \frac{\operatorname{Im} F(X')}{X' - X} \frac{1}{1 + (\lambda X')^2} dX'$$

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$$F(X | \lambda) = \frac{A}{\pi} \frac{1}{1 + (\lambda X)^2} \left[ X e^{B/X} \operatorname{Ei}(-B/X) \right. \\ \left. + \frac{1}{\lambda} \operatorname{Im}(e^{-i\lambda B} (i + \lambda X)(\pi + i \operatorname{Ei}(i\lambda B))) \right]$$

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The Cauchy integral is only predictive for high moments.

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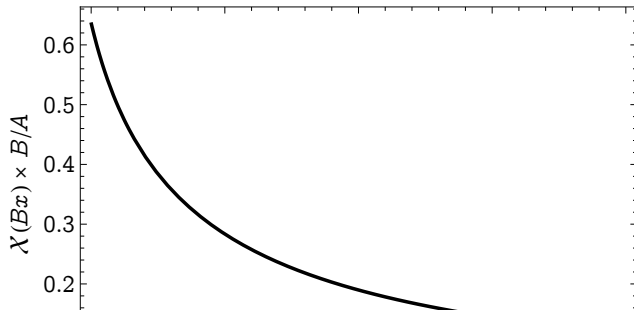
What about the susceptibility  $\chi = \frac{\partial^2 F}{\partial h^2}$ ?

Has a well-defined limit as  $\lambda \rightarrow 0$ , simple functional form:

$$\chi = t^{-\gamma} \mathcal{X}(h/t^{\beta\delta})$$

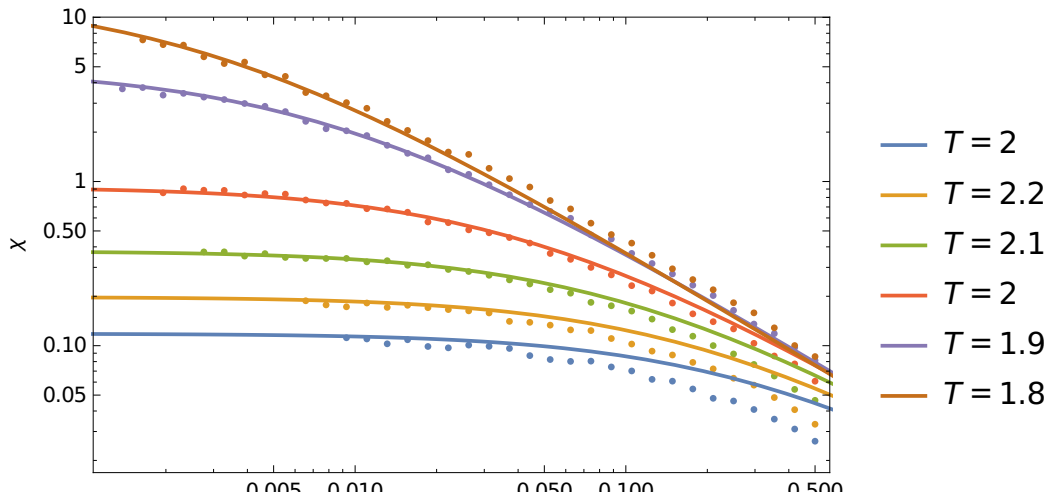
where the scaling function is

$$\mathcal{X}(X) = \frac{A}{\pi X^3} [(B - X)X + B^2 e^{B/X} \text{Ei}(-B/X)]$$



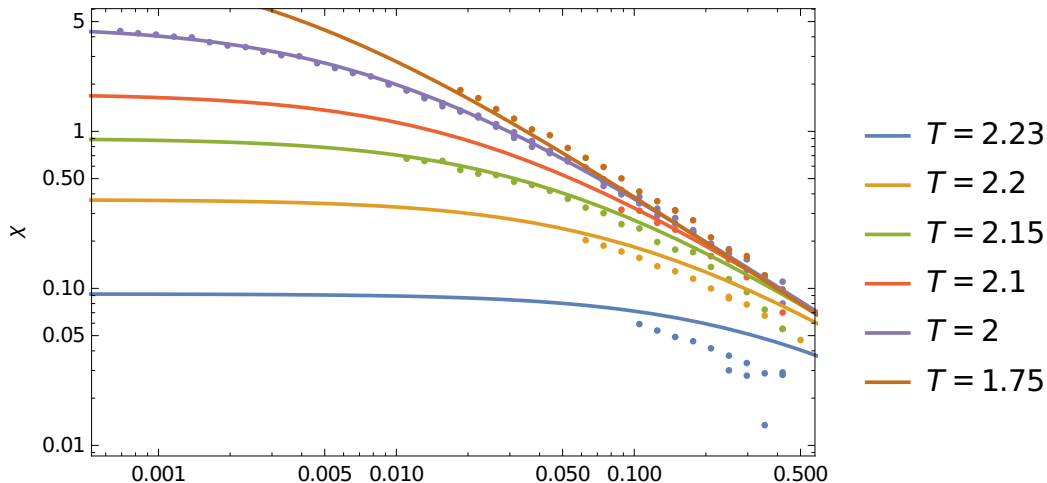
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Two parameter fit to simulations yields  $A = -0.0939(8)$ ,  $B = 5.45(6)$ , close agreement in limit of small  $t$  and  $H$ !



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Remain on the lookout for other universal properties to incorporate.

*Questions?*