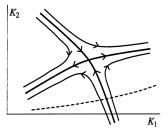
Universal scaling and the essential singularity at the abrupt Ising transition

Jaron Kent-Dobias<sup>1</sup> James Sethna<sup>1</sup>

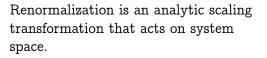
<sup>1</sup>Cornell University

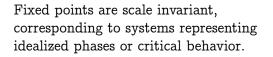
16 March 2016

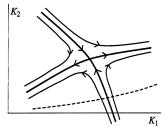
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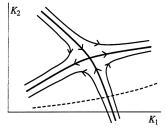
From Scaling and Renormalization in Statistical Physics by John Cardy







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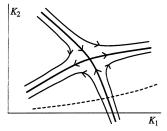


From Scaling and Renormalization in Statistical Physics by John Cardy

Renormalization is an analytic scaling transformation that acts on system space.

Fixed points are scale invariant, corresponding to systems representing idealized phases or critical behavior.

Nonanalytic behavior—like power laws and logarithms—are preserved under RG and shared by *any* system that flows to the same point.



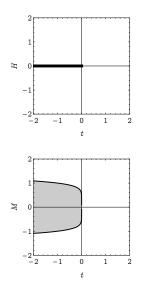
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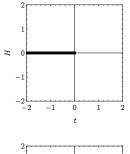
Fixed points are scale invariant, corresponding to systems representing idealized phases or critical behavior.

Nonanalytic behavior—like power laws and logarithms—are preserved under RG and shared by *any* system that flows to the same point.

Not all nonanalytic behavior is singular!



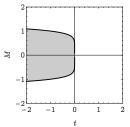
Consider an Ising-class model brought into a metastable state.

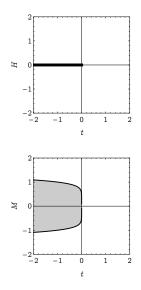


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A domain of N spins entering the stable phase causes a free energy change

 $\Delta F = \Sigma N^{\sigma} - MHN$ 





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The metastable phase is stable to domains smaller than

$$N_{
m crit} = \left(rac{MH}{\sigma\Sigma}
ight)^{-1/(\sigma-1)}$$

but those larger will grow to occupy the entire system.

The formation of a critical domain has energy cost

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Decay of the equilibrium state implies existence of an imaginary part in the free energy,

$${
m Im}\,F\sim e^{-eta\Delta F_{
m crit}}$$

Near the Ising critical point,  $\sigma = 1 - \frac{1}{d}$  and

$$M = t^{eta} \mathcal{M}(h/t^{eta \delta}) \qquad \qquad \Sigma = t^{\mu} \mathcal{S}(h/t^{eta \delta})$$

with  $\mathcal{M}(0)$  and  $\mathcal{S}(0)$  nonzero and finite.

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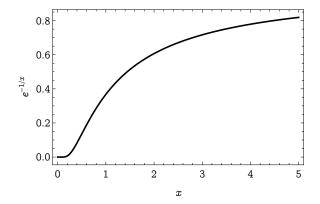
with  $\mathcal{M}(0)$  and  $\mathcal{S}(0)$  nonzero and finite.

Therefore,

$$\Delta F_{
m crit} \sim \Sigma igg( rac{MH}{\Sigma} igg)^{-(d-1)} = X^{-(d-1)} \mathcal{F}(X)$$

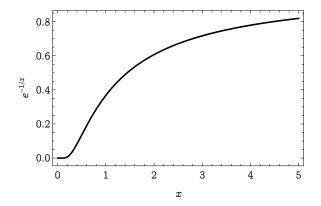
for  $X=h/t^{eta\delta}$  , and  ${
m Im}\,F={\cal I}(X)e^{-eta/X^{(d-1)}}$ 

#### The Essential Singularity



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This and its implications are therefore a universal feature of the Ising class.

Analytic properties of the partition function imply that

$$F(X) = rac{1}{\pi} \int_{-\infty}^0 rac{\operatorname{Im} F(X')}{X' - X} \, \mathrm{d} X'$$

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Only predictive for high moments of F, or

$$f_n = rac{1}{\pi} \int_{-\infty}^0 rac{\operatorname{Im} F(X')}{X'^{n+1}} \, \mathrm{d} X'$$

for  $F = \sum f_n X^n$ .

Results from field theory indicate that  $\mathcal{I}(X) \propto X$  for d = 2 and small X, so that

 $\operatorname{Im} F = AXe^{-\beta/X^{(d-1)}}$ 

$$f_n = rac{A\Gamma(n-1)}{\pi(-B)^{n-1}}$$

Not a convergent series—the real part of F for H > 0 is also nonanalytic.

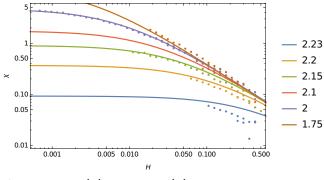
#### The Essential Singularity

In two dimensions, the Cauchy integral does not converge, normalize with

$$F(X \mid \lambda) = rac{1}{\pi} \int_{-\infty}^0 rac{\operatorname{Im} F(X')}{X' - X} rac{1}{1 + (\lambda X')^2} \, \mathrm{d} X'$$

$$\frac{A}{\pi} \frac{X e^{B/X} \operatorname{Ei}(-B/X) + \frac{1}{\lambda} \operatorname{Im}(e^{-i\lambda B}(i+\lambda X)(\pi+i\operatorname{Ei}(i\lambda B)))}{1 + (\lambda X)^2}$$

$$egin{aligned} \chi &= t^{-\gamma} \mathcal{X}(h/t^{eta\delta}) \ \mathcal{X}(X) &= rac{A}{\pi X^3} [(B-X)X + B^2 e^{B/X} \operatorname{Ei}(-B/X)] \end{aligned}$$



A = -0.0939(8), B = 5.45(6).